The Coupled Three-Body Problem and Ballistic Lunar Capture

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Efficient Lunar Transfer: How to “Shoot the Moon”

• Our goal: Design an efficient trajectory to the Moon using less fuel than the traditional Hohmann transfer (e.g., the Apollo missions).

• Find and classify sets of solutions to the Sun-Earth-Moon-spacecraft 4-body problem; in particular, find solutions starting from Earth and ending in ballistic capture by the Moon.

• Approximate 4-body system as two coupled 3-body systems.

• Review results on $L_1$ and $L_2$ dynamical channels in the planar circular restricted 3-body problem.

• Find intersections between dynamical channels between the two systems to construct complete Earth-to-Moon trajectory, which utilizes perturbation by Sun.
Lunar Capture: How to get to the Moon Cheaply

• In 1991, the failed Japanese mission, Muses-A (Uesugi [1986]), was given new life with a radical new mission concept and renamed as the **Hiton Mission** (Tanabe et al. [1982], Belbruno [1987], Belbruno and Miller [1993]).

Numerical simulation of a ballistic capture transfer trajectory for the Japanese spacecraft Hiton: ecliptic plane projection, sun's direction indicated at Earth injection. (from Belbruno and Miller [1993])
We present an approach to the problem of the orbital dynamics of this interesting trajectory by implementing in a systematic way the view that the Sun-Earth-Moon-spacecraft 4-body system can be modeled as two coupled 3-body systems.
Below is a schematic of the “Shoot the Moon” trajectory, showing the two legs of the trajectory in the Sun-Earth rotating frame:

- **Earth backward targeting portion**
- **Lunar capture portion**
- Within each 3-body system, using our understanding of the invariant manifold structures associated with the Lagrange points $L_1$ and $L_2$, we transfer from a 200 km altitude Earth orbit into the region where the invariant manifold structure of the Sun-Earth Lagrange points interact with the invariant manifold structure of the Earth-Moon Lagrange points.
- Schematic with Lagrange point invariant manifold structures:
  - **Earth backward targeting portion**
  - **Lunar capture portion**

![Diagram](image-url)
One utilizes the sensitivity of the “twisting” near the invariant manifold tubes to target back to a suitable Earth parking orbit.
This interaction permits a fuel efficient transfer from the Sun-Earth system to the Earth-Moon system. The invariant manifold tubes of the Earth-Moon system provide the dynamical channels in phase space that enable ballistic captures of the spacecraft by the Moon.
The results are then checked by integration in the bicircular 4-body problem. It works!
This technique is cheaper (about 20% less $\Delta V$) than the usual Hohmann transfer (jumping onto an ellipse that reaches the Moon, then accelerating to catch it, then circularizing). However, it also takes longer (4 to 6 months compared to 3 days).

(from Brown [1992])
Planar Circular Restricted Three Body Problem–PCR3BP

General Comments

- Describes the motion of a body moving in the gravitational field of two main bodies (the *primaries*) that are moving in circles.
- The two main bodies could be the *Sun and Earth*, or the *Earth and Moon*, etc. The total mass is normalized to 1; they are denoted \( m_S = 1 - \mu \) and \( m_E = \mu \), so \( 0 < \mu < \frac{1}{2} \).
- Let \( \mu \) be the ratio between the mass of the Earth and the mass of the Sun-Earth system,

\[
\mu = \frac{m_E}{m_E + m_S},
\]

- For the Sun-Earth system, \( \mu = 3.03591 \times 10^{-6} \).
- For the Earth-Moon system, \( \mu = 0.01215 \).
• The two main bodies rotate in the plane in circles counterclockwise about their common center of mass and with angular velocity $\omega$ (also normalized to 1).

• The third body, the *spacecraft*, has mass zero and is free to move in the plane.

• The *planar* restricted three body problem is used for simplicity. Generalization to the *three dimensional problem* is of course important, but many of the effects can be described well with the planar model.
Equations of Motion

- **Notation:** Choose a *rotating coordinate system* so that
  - the origin is at the center of mass
  - the Sun and Earth are on the $x$-axis at the points $(-\mu, 0)$ and $(1 - \mu, 0)$ respectively—i.e., the distance from the Sun to Earth is normalized to be 1.
  - Let $(x, y)$ be the position of the comet in the plane relative to the positions of the Sun and Earth
  - distances to the Sun and Earth:
    $$ r_1 = \sqrt{(x + \mu)^2 + y^2} \quad \text{and} \quad r_2 = \sqrt{(x - 1 + \mu)^2 + y^2}. $$
\[ x = -\mu \]

\[ m_S = 1 - \mu \]

\[ m_E = \mu \]

Earth’s orbit

Earth

Sun

spacecraft

\[ r_1 \]

\[ r_2 \]
• Equations of motion:

\[
\begin{align*}
\ddot{x} - 2\dot{y} &= \frac{\partial \Omega}{\partial x}, \\
\ddot{y} + 2\dot{x} &= \frac{\partial \Omega}{\partial y}
\end{align*}
\]

where

\[
\Omega = \frac{(x^2 + y^2)}{2} + \frac{1 - \mu}{r_1} + \frac{\mu}{r_2} + \frac{\mu(1 - \mu)}{2}.
\]

• They have a first integral, the Jacobi constant, given by

\[
C(x, y, \dot{x}, \dot{y}) = -(\dot{x}^2 + \dot{y}^2) + 2\Omega(x, y),
\]

which is related to the Hamiltonian energy by \( C = -2H \).

• Energy manifolds are 3-dimensional surfaces foliating the 4-dimensional phase space.

• For fixed energy, Poincaré sections are then 2-dimensional, making visualization of intersections between sets in the phase space particularly simple.
**Five Equilibrium Points**

- Three *collinear* (Euler, 1750) on the \( x \)-axis— \( L_1, L_2, L_3 \)
- Two *equilateral points* (Lagrange, 1760)— \( L_4, L_5 \).
Stability of Equilibria

- Eigenvalues of the linearized equations at $L_1$ and $L_2$ have one real and one imaginary pair. The \textit{stable and unstable manifolds} of these equilibria play an important role.

- Associated periodic orbits are called the \textit{Lyapunov orbits}. Their stable and unstable manifolds are also important.

- For space mission design, the most interesting equilibria are the \textit{unstable} ones, not the stable ones!

\textit{Consider the dynamics plus the control!}

Control often makes unstable objects, not attractors, of interest!\footnote{This is related to \textit{control of chaos}; see Bloch, A.M. and J.E. Marsden [1989] Controlling homoclinic orbits, \textit{Theor. & Comp. Fluid Mech.} 1, 179–190.} Under proper \textit{control management} they are incredibly \textit{energy efficient}. 
**Hill’s Regions**

- Our main concern is the behavior of orbits whose energy is just above that of $L_2$. Roughly, we refer to this small energy range as the *temporary capture energy* range.
  - Fortunately, the temporary capture energy surfaces for the Sun-Earth and Earth-Moon systems intersect in phase space, making a fuel efficient transfer possible.

- The **Hill’s region** is the projection of this energy region onto position space.

- The region not accessible for these energies is called the *forbidden region*.

- For this case, the Hill’s region contains a “neck” about $L_1$ and $L_2$. This equilibrium neck region and its relation to the global orbit structure is critical: it was studied in detail by Conley, McGehee and the Barcelona group.
Orbits with energy just above that of $L_2$ can be transit orbits, passing through the neck region between the exterior region (outside Earth’s orbit) and the Earth temporary capture region (bubble surrounding Earth). They can also be nontransit orbits or asymptotic orbits.
Equilibrium Region near a Lagrange Point

- **4 types** of orbits in the *equilibrium region*.
  - Black circle is the unstable periodic Lyapunov orbit.
  - 4 cylinders of asymptotic orbits form pieces of stable and unstable manifolds. They intersect the bounding spheres at asymptotic circles, separating *spherical caps*, which contain transit orbits, from *spherical zones*, which contain nontransit orbits.
Roughly speaking, the equilibrium region has the dynamics of a \textit{saddle \times harmonic oscillator}.
Tubes Partition the Energy Surface

- Stable and unstable manifold tubes act as separatrices for the flow in the equilibrium region.
  - Those inside the tubes are transit orbits.
  - Those outside the tubes are nontransit orbits.
  - e.g., transit from outside Earth’s orbit to Earth capture region possible only through $L_2$ periodic orbit stable tube.
- **Stable** and **unstable** manifold tubes effect the *transport* of material to and from the capture region.
• Tubes of transit orbits can be utilized for \textit{ballistic capture}.  

![Diagram showing transit orbits and trajectory]

- Trajectory Begins Inside Tube
- Ends in Ballistic Capture
- Passes through $L_2$ Equilibrium Region
Invariant manifold tubes are *global objects* — extend far beyond vicinity of Lagrange points.
Twisting of Orbits in the Equilibrium Region

- Recall that the equilibrium region roughly has the dynamics of a saddle $\times$ harmonic oscillator.
  - Orbits *twist* in the equilibrium region, roughly following the Lyapunov orbit.
  - The closer the approach to the Lyapunov orbit, the more the orbit twists.
  - Thus, the closer an orbit begins to the tube on its approach to the equilibrium region, the more it will be twisted when it exits the equilibrium region.

- This twisting near the tubes will be used in the Earth backward targeting portion of the Earth-to-Moon trajectory.
Twisting in the equilibrium region can be understood in terms of mappings between the bounding spheres.
- We select a line of constant $x$-position passing through the Earth in the Earth capture region.
- Pick a few points near the **unstable tube**, with the same position but slightly different velocities.
• With a slight change in the velocity, we can target any position on the lower line, where the mirror image stable tube intersects.
• Look at Poincaré section along this line of constant $x$-position.

• With infinitesimal changes in velocity, any point near lower tube cross-section can be targeted.

![Poincare Section Diagram](image-url)

**Earth Targeting Using "Twisting"**

- **Pre-Image of Strip $S$**
- **Stable Manifold**
- **Unstable Manifold**
- **Earth**
- **Sun**
- **$L_2$ orbit**
Two Coupled 3-Body Systems

- We obtain the fuel efficient transfer by taking full advantage of the dynamics of the 4-body system (Earth, Moon, Sun, and spacecraft) by initially modeling it as two coupled planar circular restricted 3-body systems.

- In this approach, we utilize the Lagrange point dynamics of both the Earth-Moon-spacecraft and Sun-Earth-spacecraft systems.
In this simplified model, the Moon is on a circular orbit about the Earth, the Earth (or rather the Earth-Moon center of mass) is on a circular orbit about the Sun, and the systems are coplanar.

In the actual solar system:
- Moon’s eccentricity is 0.055
- Earth’s eccentricity is 0.017
- Moon’s orbit is inclined to Earth’s orbit by 5°

These values are low, so the coupled planar circular 3-body problem is considered a good starting model.

An orbit which becomes a real mission is typically obtained first in such an approximate system and then later refined through more precise models which include effects such as out-of-plane motion, eccentricity, the other planets, solar wind, etc.

However, tremendous insight is gained by considering a simpler model which reveals the essence of the transfer dynamics.
Notice that in the Sun-Earth rotating frame, patterns such as the Sun-Earth $L_2$ portion of the trajectory are made plain which are not discernable in the inertial frame.
The construction is done mainly in the Sun-Earth rotating frame using the Poincaré section (that passes through the $x$-position of the Earth) which help to glue the Sun-Earth Lagrange point orbit portion of the trajectory with the lunar ballistic capture portion.
Construction of Earth-to-Moon Transfer

- **Strategy**: Find an initial condition (position and velocity) for a spacecraft on the Poincaré section such that:
  - when integrating *forward*, spacecraft will be guided by *Earth-Moon manifold* and get ballistically captured by the Moon;
  - when integrating *backward*, spacecraft will hug *Sun-Earth manifolds* and return to Earth.

- We utilize two important properties of Lagrange point dynamics:
  - Tube of *transit orbits* is key in finding capture orbit for Earth-Moon portion of the design;
  - *Twisting* of orbits in equilibrium region is key in finding a fuel efficient transfer for Sun-Earth portion of the design.
Lunar Ballistic Capture Portion

- Recall that by targeting the region enclosed by the stable manifold tube of the Earth-Moon-S/C system’s $L_2$ Lyapunov orbit, we can construct an orbit which will get ballistically captured by the Moon.
• When we transform the Poincaré section of the stable manifold of the Lyapunov orbit about the Earth-Moon $L_2$ point into the Sun-Earth rotating frame, we obtain a curve. A point interior to this curve, with the correct phasing of the Moon, will approach the Moon when integrated forward. The phasing of the Moon is determined by the orientation of the Poincaré section in the Earth-Moon system.
• Assuming the Sun is a negligible perturbation to the Earth-Moon-S/C 3-body dynamics, any spacecraft with initial conditions within this **closed loop** will be ballistically captured by the Moon.
• “Ballistic capture by the Moon” means an orbit which under natural dynamics gets within the sphere of influence of the Moon (20,000 km) and performs at least one revolution around the Moon. In such a state, a slight $\Delta V$ will result in a stable capture (closing off the necks at $L_1$ and $L_2$).
**Earth Backward Targeting Portion**

- Pick an energy in the temporary capture range of the Sun-Earth system which has $L_2$ orbit manifolds that come near a 200 km altitude Earth parking orbit.
• Compute Poincaré section along line of constant $x$-position passing through Earth:

  - The **red curve** is the Poincaré cut of the **unstable manifold** of the Lyapunov orbit around the Sun-Earth $L_2$. 

![Diagram showing Poincaré section and orbit dynamics](image)
Picking an initial condition just outside this curve, we can backward integrate to produce a trajectory coming back to the Earth parking orbit.
Connecting the Two Portions

- Vary the phase of the Moon until Earth-Moon $L_2$ manifold curve intersects Sun-Earth $L_2$ manifold curve.
• Intersection is found!

• In the region which is in the interior of the green curve but in the exterior of the red curves,
• an orbit will get ballistically captured by the Moon when integrated forward;
• when integrated backward, orbit will hug the **unstable** manifold back to the Sun-Earth $L_2$ equilibrium region with a twist, and then hug the **stable** manifold back towards Earth.
• With only a slight modification, a midcourse $\Delta V$ at the patch point (34 m/s), this procedure produces a genuine solution integrated in the bicircular 4-body problem.

• Since capture at Moon is natural (zero $\Delta V$), the amount of on-board fuel necessary is lowered (by about 20%).
Why Does It Work?

- Heuristic arguments for using the coupled 3-body model:
  - When outside the Moon’s small sphere of influence (20,000 km), which is most of pre-capture flight, we can consider the Moon’s perturbation on the Sun-Earth-S/C 3-body system to be negligible. Thus, can utilize Sun-Earth Lagrange point invariant manifold structures.
  - The midcourse $\Delta V$ is performed at a point where the spacecraft is entering the Earth’s sphere of influence (900,000 km), where we can consider the Sun’s perturbation on the Earth-Moon-S/C 3-body system to be negligible. Thus, Earth-Moon Lagrange point structures can be utilized for capture sequence.
What Next? — Future Work

- Using differential correction and continuation, we expect this trajectory can be used as an initial guess to generalize to the three-dimensional 4-body problem as well as the full solar system model.

- Optimize trajectory (minimize total fuel consumption):
  - Incorporate initial lunar swingby
  - Apply optimal control (e.g., COOPT)
  - Use continuous thrust (low-thrust)

- Develop systematic procedure for coupling multiple 3-body systems. This will aid in the design of innovative space missions and help aid in understanding subtle non-Keplerian transport throughout the solar system.
More Information and References


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